

# Effective gluon mass and freezing of the QCD coupling

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Based on:

A. C. Aguilar and J. Papavassiliou, *In preparation*

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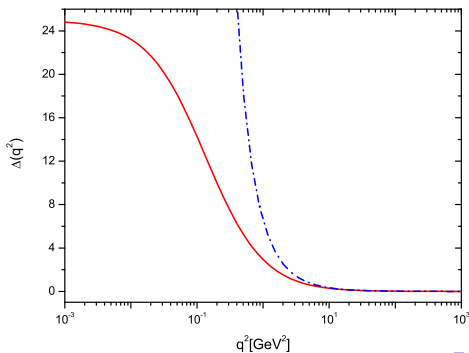
“Gluon mass generation in the *PT*-BFM scheme,”  
*JHEP* **0612**, 012 (2006) [arXiv:hep-ph/0610040]

# General Considerations

Gluon “propagator”  $\Delta_{\mu\nu}(q)$

$$i\Delta_{\mu\nu}(q) = \left[ g_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right] \Delta(q^2) + \frac{q_\mu q_\nu}{q^4}.$$

**Dynamical generation** of an **infrared cutoff**. The QCD dynamics allow for  $\Delta^{-1}(0) \neq 0$ . Acts as an **effective “mass”** for the gluons. **Cornwall, Phys. Rev. D 26, 1453 (1982)**



- Does **not** correspond to a term  $m^2 A_\mu^2$  in the QCD Lagrangian.
- The **local gauge symmetry remains exact**.
- **Not hard** but **momentum dependent** mass:  $m = m(q^2)$ .
- **Drops off** in the UV “sufficiently fast”.  $\implies$  QCD remains **renormalizable**
- Purely **non-perturbative effect**:
  - **Lattice** (discretized space-time)
  - **Schwinger-Dyson** equations (continuum).

# Schwinger-Dyson Equation

$$\Delta^{-1}(q^2)P_{\mu\nu}(q) = q^2P_{\mu\nu}(q) + i [\Pi_{\mu\nu}^{(a_1)}(q) + \Pi_{\mu\nu}^{(a_2)}]$$

$$\Pi_{\mu\nu}^{(a_1)}(q) = \frac{1}{2} C_A g^2 \int [dk] \Gamma_{\mu\alpha\beta} \Delta^{\alpha\alpha'}(k) \tilde{\Gamma}_{\nu\alpha'\beta'} \Delta^{\beta\beta'}(k+q)$$

$$\Pi_{\mu\nu}^{(a_2)} = -C_A g^2 g_{\mu\nu} \int [dk] \Delta(k)$$

The expression for the vertex that we will use is given by

$$\tilde{\Gamma}^{\mu\alpha\beta} = L^{\mu\alpha\beta} + T_1^{\mu\alpha\beta} + T_2^{\mu\alpha\beta}$$

with

$$L^{\mu\alpha\beta}(q, p_1, p_2) = \tilde{\Gamma}^{\mu\alpha\beta}(q, p_1, p_2) + ig^{\alpha\beta} \frac{q^\mu}{q^2} [\Pi(p_2) - \Pi(p_1)]$$

$$T_1^{\mu\alpha\beta}(q, p_1, p_2) = -i \frac{c_1}{q^2} (q^\beta g^{\mu\alpha} - q^\alpha g^{\mu\beta}) [\Pi(p_1) + \Pi(p_2)]$$

$$T_2^{\mu\alpha\beta}(q, p_1, p_2) = -ic_2 (q^\beta g^{\mu\alpha} - q^\alpha g^{\mu\beta}) \left[ \frac{\Pi(p_1)}{p_1^2} + \frac{\Pi(p_2)}{p_2^2} \right]$$

$$\tilde{\Gamma}_{\mu\alpha\beta}(q, p_1, p_2) = (p_1 - p_2)_\mu g_{\alpha\beta} + 2q_\beta g_{\mu\alpha} - 2q_\alpha g_{\mu\beta}$$

$$\Delta^{-1}(x) = Kx + bg^2 \sum_{i=1}^8 a_i A_i(x) + \Delta^{-1}(0)$$

$$A_1(x) = a_1 x \int_x^\infty dy y \Delta^2(y)$$

$$A_2(x) = a_2 x \int_x^\infty dy \Delta(y)$$

$$A_3(x) = a_3 x \Delta(x) \int_0^x dy y \Delta(y)$$

$$A_4(x) = a_4 \int_0^x dy y^2 \Delta^2(y)$$

$$A_5(x) = a_5 \Delta(x) \int_0^x dy y^2 \Delta(y)$$

$$A_6(x) = a_6 \int_0^x dy y \Delta(y)$$

$$A_7(x) = a_7 \frac{\Delta(x)}{x} \int_0^x dy y^3 \Delta(y)$$

$$A_8(x) = a_8 \frac{1}{x} \int_0^x dy y^3 \Delta^2(y)$$

The renormalization condition  $K$  is fixed by  $\Delta^{-1}(\mu^2) = \mu^2$

# The UV behavior of effective gluon mass

$$m^2(x) \ln x = g^{-2} \Delta^{-1}(0) + \gamma_1 \int_0^x dy m^2(y) \tilde{\Delta}(y) + \frac{\gamma_2}{x} \int_0^x dy y m^2(y) \tilde{\Delta}(y)$$

with

$$\tilde{\Delta}(q^2) = \frac{1}{q^2 + m^2(q^2)},$$

The asymptotic solutions:

$$m^2(x) = \frac{\lambda_2^4}{x} (\ln x)^{\gamma_2 - 1} \implies \langle G_{\mu\nu}^a G_a^{\mu\nu} \rangle$$

# Propagator and Running Masses

The RG quantity,  $d(q^2) = g^2 \Delta(q^2)$ , has the general form:

$$d(q^2) = \frac{\bar{g}^2(q^2)}{q^2 + m^2(q^2)},$$

where the **dynamical mass**  $m^2(q^2)$  and **effective charge**  $\bar{g}^2(q^2)$  are

$$m^2(q^2) = \frac{m_0^4}{q^2 + m_0^2} \left[ \ln \left( \frac{q^2 + \rho m_0^2}{\Lambda^2} \right) / \ln \left( \frac{\rho m_0^2}{\Lambda^2} \right) \right]^{\gamma_2 - 1}$$

$$\bar{g}^2(q^2) = \left[ b \ln \left( \frac{q^2 + \rho m^2(q^2)}{\Lambda^2} \right) \right]^{-1}$$



# Numerical Results

