

Classification of the Flipped $SU(5)$ Heterotic-String Vacua

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- Heterotic String Phenomenology
- Free Fermionic Construction
- Flipped $SU(5)$ Models
- Classification
- Future Work

Free Fermionic Construction

- 4D Theory
- $N = 1$ Supersymmetry
- 3 Generation Standard Model Fermions
- $SO(10)$ GUTs
- Absence of exotic states

Free Fermionic Construction

Properties

- Conformally invariance
- Decoupling left and right moving modes
- $D = 4$ theory

Result

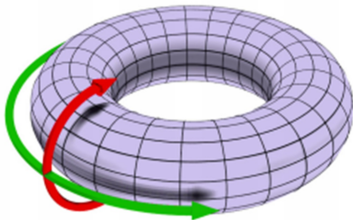
- $C_L = -26 + 11 + D + \frac{D}{2} + \frac{N_{f_L}}{2} = 0$
 \implies 18 left-moving real fermions
- $C_R = 0$
 \implies 44 right-moving real fermions

Free Fermionic Construction

- Partition function is used to include all physical states

$$Z = \sum_{\alpha, \beta} c \begin{pmatrix} \alpha \\ \beta \end{pmatrix} Z [\alpha, \beta]$$

- Taking the one-loop partition function transforms the worldsheet into a torus.



Free Fermionic Construction

$$\alpha = \left\{ \psi^{1,2}, \chi^i, y^i, \omega^i | \bar{y}^i, \bar{\omega}^i, \bar{\psi}^{1,\dots,5}, \bar{\eta}^{1,2,3}, \bar{\phi}^{1,\dots,8} \right\}$$

Where $i = 1, \dots, 6$

- Left-movers

- X_L^μ , $\mu = 1, 2$ 2 transverse coordinates
- ψ_L^μ , $\mu = 1, 2$ The fermionic partners
- Ω^j , $j = 1, \dots, 18$ 18 internal real fermions

- Right-movers

- X_R^μ , $\mu = 1, 2$ 2 transverse coordinates
- $\bar{\Omega}^j$, $j = 1, \dots, 44$ 44 internal real fermions

Free Fermionic Construction

- ABK Rules

- $\sum_i m_i b_i = 0$
- $N_{ij} \cdot b_i \cdot b_j = \text{mod } 4$
- $N_i \cdot b_i \cdot b_i = \text{mod } 8$
- $1 \in \Xi$
- Even number of fermions

- One-Loop Phases

- $C \begin{pmatrix} b_i \\ b_j \end{pmatrix} = \pm 1 \text{ or } \pm i$

- GSO Projection

- $e^{i\pi b_i \cdot F_\alpha} |s\rangle_\alpha = \delta_\alpha C \begin{pmatrix} \alpha \\ b_i \end{pmatrix}^* |s\rangle_\alpha$

- Virasoro Condition

- $M_L^2 = -\frac{1}{2} + \frac{\alpha_L^2}{8} + \sum v_L = -1 + \frac{\alpha_R^2}{8} + \sum v_R = M_R^2$

Flipped $SU(5)$ Basis Vectors

Basis Vectors

- $v_1 = S = \{\psi^\mu, \chi^{1,\dots,6}\}$
- $v_{1+i} = e_i = \{y^i, \omega^i | \bar{y}^i, \bar{\omega}^i\} \quad i = 1, \dots, 6$
- $v_8 = b_1 = \{\chi^{34}, \chi^{56}, y^{34}, y^{56} | \bar{y}^{34}, \bar{y}^{56}, \bar{\eta}^1, \bar{\psi}^{1,\dots,5}\}$
- $v_9 = b_2 = \{\chi^{12}, \chi^{56}, y^{12}, y^{56} | \bar{y}^{12}, \bar{y}^{56}, \bar{\eta}^2, \bar{\psi}^{1,\dots,5}\}$
- $v_{10} = z_1 = \{\bar{\phi}^{1,\dots,4}\}$
- $v_{11} = z_2 = \{\bar{\phi}^{5,\dots,8}\}$
- $v_{12} = \alpha = \{\bar{\eta}^{1,2,3} = \frac{1}{2}, \bar{\psi}^{1,\dots,5} = \frac{1}{2}, \bar{\phi}^{1,\dots,4} = \frac{1}{2}, \bar{\phi}^5 = 1\}$

Gauge Group

$$SU(5) \times U(1)_5 \times U(1)_1 \times U(1)_2 \times U(1)_3 \times \textit{Hidden}$$

Depending on the choices of the projection coefficients, extra gauge bosons may arise. These arise with any linear combination of z_1 , z_2 and α , which are massless. Such as

$$\mathbf{G} = \left\{ \begin{array}{cccc} z_1, & z_2, & z_1 + z_2, & \alpha, \\ \alpha + z_1, & \alpha + z_2, & \alpha + z_1 + z_2, & x \end{array} \right\}$$

Where

$$x = 2\alpha + z_1 = \{\bar{\psi}^{1,\dots,5}, \bar{\eta}^{1,2,3}\}$$

Observable Gauge Group Enhancements

x is a sector which can enlarge the observable gauge group. Enhancement takes place when the following conditions are satisfied

Sector Condition	
$(z_1 + 2\alpha e_i) = (z_1 + 2\alpha z_k) = 0$	
Enhancement Condition	Resulting Enhancement
$(z_1 + 2\alpha \alpha) = (z_1 + 2\alpha b_2)$	$SU(5)_{obs} \times U(1)_5 \times U(1)_1 \times U(1)_2 \times U(1)_3$ $\longrightarrow SU(6) \times SU(2) \times U(1)^2$
$(z_1 + 2\alpha \alpha) \neq (z_1 + 2\alpha b_2)$	$SU(5)_{obs} \times U(1)_5 \times U(1)_1 \times U(1)_2 \times U(1)_3$ $\longrightarrow SO(10) \times U(1)^3$

The pre-stated conditions hold for all $i = 1, \dots, 6$.

GGSO Coefficients

$$(v_i|v_j) =$$

	S	e ₁	e ₂	e ₃	e ₄	e ₅	e ₆	b ₁	b ₂	z ₁	z ₂	α
S	1	1	1	1	1	1	1	1	1	1	1	1
e ₁	1	0	0	1	1	1	0	0	0	1	1	1
e ₂	1	0	0	1	1	1	0	0	0	1	1	0
e ₃	1	1	1	0	0	1	0	0	0	1	1	0
e ₄	1	1	1	0	0	0	0	0	1	0	0	1
e ₅	1	1	1	1	0	0	0	0	1	0	1	0
e ₆	1	0	0	0	0	0	0	1	1	0	0	1
b ₁	0	0	0	0	0	0	1	1	0	1	1	1/2
b ₂	0	0	0	0	1	1	1	0	1	0	1	-1/2
z ₁	1	1	1	1	0	0	0	1	0	0	0	0
z ₂	1	1	1	1	0	1	0	1	1	0	0	1/2
α	1	1	0	0	1	0	1	0	1	1	1	1

Where $C\left(\begin{smallmatrix} v_i \\ v_j \end{smallmatrix}\right) = e^{i\pi(v_i|v_j)}$ $(b_i|b_j) \in \{-\frac{1}{2}, 0, \frac{1}{2}, 1\}$

GGSO Coefficients

$$\begin{array}{c}
 S \\
 e_1 \\
 e_2 \\
 e_3 \\
 e_4 \\
 e_5 \\
 e_6 \\
 b_1 \\
 b_2 \\
 z_1 \\
 z_2 \\
 \alpha
 \end{array}
 \begin{pmatrix}
 S & e_1 & e_2 & e_3 & e_4 & e_5 & e_6 & b_1 & b_2 & z_1 & z_2 & \alpha \\
 & & 0_1 & 1_2 & 1_3 & 1_4 & 0_5 & 0_6 & & 1_7 & & 1_8 \\
 & & & 1_9 & 1_{10} & 1_{11} & 0_{12} & 0_{13} & & 1_{14} & & 0_{15} \\
 & & & & 0_{16} & 1_{17} & 0_{18} & & 0_{19} & 1_{20} & & 0_{21} \\
 & & & & & 0_{22} & 0_{23} & & 1_{24} & 0_{25} & & 1_{26} \\
 & & & & & & 0_{27} & 0_{28} & 1_{29} & 0_{30} & & 0_{31} \\
 & & & & & & & 1_{32} & 1_{33} & 0_{34} & & 1_{35} \\
 & & & & & & & & 0_{36} & & 1_{37} & 1/2_{38} \\
 & & & & & & & & & 0_{39} & 1_{40} & -1/2_{41} \\
 & & & & & & & & & & 0_{42} & 0_{43} \\
 & & & & & & & & & & & 1/2_{44}
 \end{pmatrix}$$

The chiral matter spectrum arises from the twisted sectors. The chiral spinorial representations of the observable $SU(5) \times U(1)$ arise from the sectors

$$\begin{aligned} B_{pqrs}^{(1)} &= S + b_1 + pe_3 + qe_4 + re_5 + se_6, \\ &= \{\psi^\mu, \chi^{12}, (1-p)y^3\bar{y}^3, p\omega^3\bar{\omega}^3, (1-q)y^4\bar{y}^4, q\omega^4\bar{\omega}^4, \\ &\quad (1-r)y^5\bar{y}^5, r\omega^5\bar{\omega}^5, (1-s)y^6\bar{y}^6, s\omega^6\bar{\omega}^6, \bar{\eta}^1, \bar{\psi}^{1,\dots,5}\}, \end{aligned}$$

$$B_{pqrs}^{(2)} = S + b_2 + pe_1 + qe_2 + re_5 + se_6,$$

$$B_{pqrs}^{(3)} = S + b_3 + pe_1 + qe_2 + re_3 + se_4,$$

Where $p, q, r, s = 0, 1$ and $b_3 = b_1 + b_2 + 2\alpha + z_1$.

The states in the sector $B_{pqrs}^{(1)}$ can be projected out of the spectrum by the GGSO projection of the vectors e_1 , e_2 , z_1 and z_2 . Similarly for all sectors, we can define a projector P such that the states survive when $P = 1$ and are projected out when $P = 0$:

$$P_{pqrs}^{(1)} = \frac{1}{16} \left(1 - C \left(\begin{matrix} e_1 \\ B_{pqrs}^{(1)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} e_2 \\ B_{pqrs}^{(1)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_1 \\ B_{pqrs}^{(1)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_2 \\ B_{pqrs}^{(1)} \end{matrix} \right) \right)$$

$$P_{pqrs}^{(2)} = \frac{1}{16} \left(1 - C \left(\begin{matrix} e_3 \\ B_{pqrs}^{(2)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} e_4 \\ B_{pqrs}^{(2)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_1 \\ B_{pqrs}^{(2)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_2 \\ B_{pqrs}^{(2)} \end{matrix} \right) \right)$$

$$P_{pqrs}^{(3)} = \frac{1}{16} \left(1 - C \left(\begin{matrix} e_5 \\ B_{pqrs}^{(3)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} e_6 \\ B_{pqrs}^{(3)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_1 \\ B_{pqrs}^{(3)} \end{matrix} \right) \right) \cdot \left(1 - C \left(\begin{matrix} z_2 \\ B_{pqrs}^{(3)} \end{matrix} \right) \right)$$

Projectors

These projectors can be expressed as a system of linear equations with p , q , r and s as unknowns. The solutions of a specific system of equations yield the different combinations of p , q , r and s for which sectors survive the GSO projections.

$$\begin{pmatrix} (e_1|e_3) & (e_1|e_4) & (e_1|e_5) & (e_1|e_6) \\ (e_2|e_3) & (e_2|e_4) & (e_2|e_5) & (e_2|e_6) \\ (z_1|e_3) & (z_1|e_4) & (z_1|e_5) & (z_1|e_6) \\ (z_2|e_3) & (z_2|e_4) & (z_2|e_5) & (z_2|e_6) \end{pmatrix} \begin{pmatrix} p \\ q \\ r \\ s \end{pmatrix} = \begin{pmatrix} (e_1|b_1) \\ (e_2|b_1) \\ (z_1|b_1) \\ (z_2|b_1) \end{pmatrix},$$

$$\begin{pmatrix} (e_3|e_1) & (e_3|e_2) & (e_3|e_5) & (e_3|e_6) \\ (e_4|e_1) & (e_4|e_2) & (e_4|e_5) & (e_4|e_6) \\ (z_1|e_1) & (z_1|e_2) & (z_1|e_5) & (z_1|e_6) \\ (z_2|e_1) & (z_2|e_2) & (z_2|e_5) & (z_2|e_6) \end{pmatrix} \begin{pmatrix} p \\ q \\ r \\ s \end{pmatrix} = \begin{pmatrix} (e_3|b_2) \\ (e_4|b_2) \\ (z_1|b_2) \\ (z_2|b_2) \end{pmatrix},$$

$$\begin{pmatrix} (e_5|e_1) & (e_5|e_2) & (e_5|e_3) & (e_5|e_4) \\ (e_6|e_1) & (e_6|e_2) & (e_6|e_3) & (e_6|e_4) \\ (z_1|e_1) & (z_1|e_2) & (z_1|e_3) & (z_1|e_4) \\ (z_2|e_1) & (z_2|e_2) & (z_2|e_3) & (z_2|e_4) \end{pmatrix} \begin{pmatrix} p \\ q \\ r \\ s \end{pmatrix} = \begin{pmatrix} (e_5|b_3) \\ (e_6|b_3) \\ (z_1|b_3) \\ (z_2|b_3) \end{pmatrix}.$$

These 48 sectors give rise to $\mathbf{16}$ and $\overline{\mathbf{16}}$ representations of $SO(10)$ decomposed under $SU(5) \times U(1)$

$$\mathbf{16} = (\overline{\mathbf{5}}, -\frac{3}{2}) + (\mathbf{10}, +\frac{1}{2}) + (\mathbf{1}, +\frac{5}{2}),$$

$$\overline{\mathbf{16}} = (\mathbf{5}, +\frac{3}{2}) + (\overline{\mathbf{10}}, -\frac{1}{2}) + (\mathbf{1}, -\frac{5}{2}).$$

The Standard Model Particles

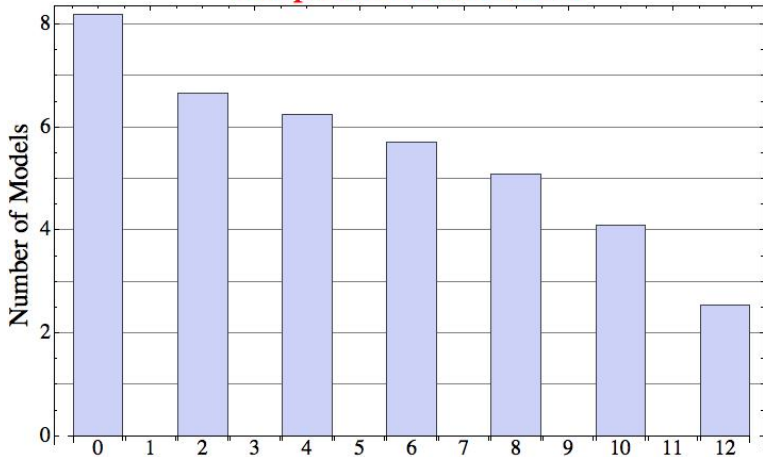
We can decompose the flipped $SU(5)$ representation under $SU(3) \times SU(2) \times U(1)$ as

$$\begin{aligned}(\bar{\mathbf{5}}, -\frac{3}{2}) &= (\bar{\mathbf{3}}, 1, -\frac{2}{3})_{u^c} + (1, 2, -\frac{1}{2})_L, \\(\mathbf{10}, +\frac{1}{2}) &= (3, 2, +\frac{1}{6})_Q + (\bar{\mathbf{3}}, 1, +\frac{1}{3})_{d^c} + (1, 1, 0)_{\nu^c}, \\(\mathbf{1}, +\frac{5}{2}) &= (1, 1, +1)_{e^c},\end{aligned}$$

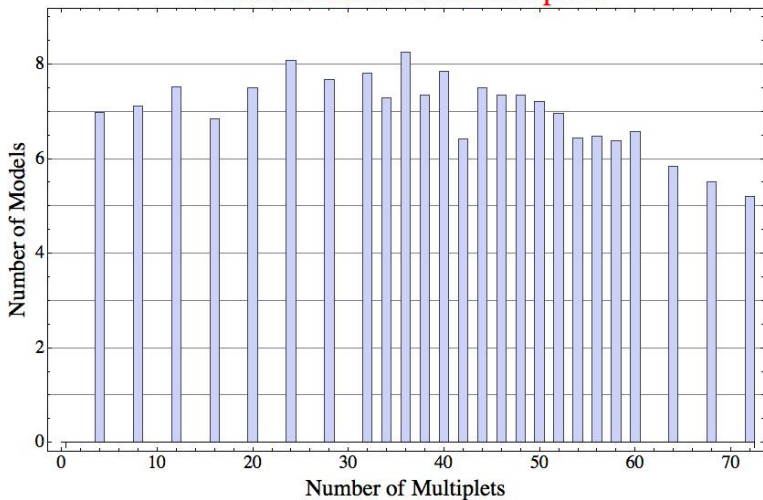
Classification

- There are 44 independent coefficients which corresponds to $2^{44} \approx 10^{13}$ different vacua.
- Using computer code we perform a statistical sampling in this space of models and extract 10^{12} distinct configurations with the flipped $SU(5)$ gauge group.

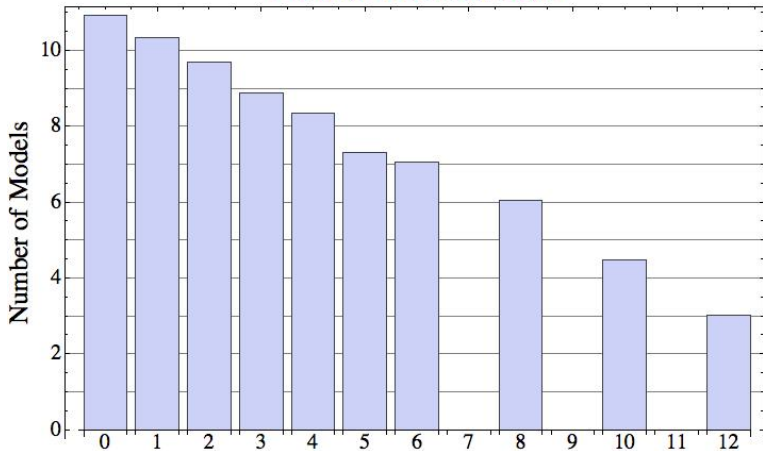
Exophobic Generations



3 Generation Exotic Multiplets



Exotic Generations



Classification

	Constraints	Total models in sample	Probability	Estimated number of models in class
	No Constraints	1000000000000	1	1.76×10^{13}
(1)	+ No Enhancements	762269298719	7.62×10^{-1}	1.34×10^{13}
(2)	+ Anomaly Free Flipped $SU(5)$	139544182312	1.40×10^{-1}	2.45×10^{12}
(3)	+ 3 Generations	738045321	7.38×10^{-4}	1.30×10^{10}
(4a)	+ SM Light Higgs	706396035	7.06×10^{-4}	1.24×10^{10}
(4b)	+ Flipped $SU(5)$ Heavy Higgs	46470138	4.65×10^{-5}	8.18×10^8
(5)	+ SM Light Higgs + & Heavy Higgs	43624911	4.36×10^{-5}	7.67×10^8
(6a)	+ Minimal Flipped $SU(5)$ Heavy Higgs	42310396	4.23×10^{-5}	7.44×10^8
(6b)	+ Minimal SM Light Higgs	25333216	2.53×10^{-5}	4.46×10^8
(7)	+ Minimal Flipped $SU(5)$ Heavy Higgs + & Minimal SM Light Higgs	24636896	2.46×10^{-5}	4.33×10^8
(8)	+ Minimal Exotic States	1218684	1.22×10^{-6}	2.14×10^7

Future work

- Extensions on Pati-Salam models
- $SU(6) \times SU(2)$
- $SU(3)_C \times U(1)_C \times SU(2)_L \times U(1)_L$
- Other Gauge Groups
- Compute the Superpotentials

THANK YOU